

The asymptotic distributions of the size of the largest length of a reversible Markov process of polymerization

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2003 J. Phys. A: Math. Gen. 36 7485

(<http://iopscience.iop.org/0305-4470/36/27/302>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.86

The article was downloaded on 02/06/2010 at 16:20

Please note that [terms and conditions apply](#).

The asymptotic distributions of the size of the largest length of a reversible Markov process of polymerization

Dong Han

Department of Mathematics, Shanghai Jiao Tong University, Shanghai, 200 030, People's Republic of China

Received 14 January 2003, in final form 28 May 2003

Published 25 June 2003

Online at stacks.iop.org/JPhysA/36/7485

Abstract

In this paper, the subcritical, near-critical and supercritical asymptotic behaviour of a reversible Markov process as a chemical model for polymerization was studied. We establish the existence of three distinct stages (subcritical, near-critical and supercritical stages) of polymerization (in the thermodynamic limit as $N \rightarrow +\infty$), depending on the value of strength of the fragmentation reaction. These three stages correspond to the size of the largest length of polymers of size N to be itself of order $\log N$, N^a ($1/2 < a < 1$) and N , respectively. Especially, $a = \frac{2}{2\sigma+1}$ when the coagulation (reaction) rate constant R_{ij} of the polymers satisfies $R_{ij} = i^\sigma j^\sigma$ with $1/2 < \sigma < 3/2$.

PACS numbers: 82.35.Jk, 02.50.Ga, 82.70.Gg

1. Introduction

A necessary and sufficient condition for the occurrence of a gelation in a reversible Markov process as a chemical model of reversible polymerization has been given in [12]. In this paper we continue to study the asymptotic behaviour of the reversible Markov process.

The process considered in the paper is a continuous-time reversible Markov chain $\{M_N(t) : t \geq 0\}$ with the state space

$$\Omega_N = \left\{ \underline{n} \in N^N : \sum_{k=1}^N kn_k = N \right\}. \quad (1)$$

The k th component of the state vector \underline{n} represents the number of k -mers. The only allowed transitions from \underline{n} are to states of the form

$$\underline{n}_{ik}^+ = \begin{cases} (n_1, n_2, \dots, n_i - 1, \dots, n_k - 1, \dots, n_{i+k} + 1, \dots, n_N) & \text{if } i \neq k \\ (n_1, n_2, \dots, n_i - 2, \dots, n_{2i} + 1, \dots, n_N) & \text{if } i = k \end{cases}$$

$$\underline{n}_{ik}^- = \begin{cases} (n_1, n_2, \dots, n_i + 1, \dots, n_k + 1, \dots, n_{i+k} - 1, \dots, n_N) & \text{if } i \neq k \\ (n_1, n_2, \dots, n_i + 2, \dots, n_{2i} - 1, \dots, n_N) & \text{if } i = k \end{cases}$$

and they occur with rate

$$Q_{\underline{n} \underline{n}'} = \begin{cases} \frac{1}{N^2} R_{ij} n_i n_j & \text{if } \underline{n}' = \underline{n}_{ij}^+ \quad i \neq j \\ \frac{1}{N^2} R_{jj} n_j (n_j - 1) & \text{if } \underline{n}' = \underline{n}_{ij}^+ \quad i = j \\ \frac{1}{N} F_{ij} n_{i+j} & \text{if } \underline{n}' = \underline{n}_{ij}^- \\ 0 & \text{other } \underline{n}' \neq \underline{n} \end{cases}$$

where the coagulation rate constants R_{ij} and fragmentation rate constants F_{ij} satisfy the following detailed balance condition proposed by Van Dongen and Ernst [25]

$$\lambda F_{ij} f(i+j) = R_{ij} f(i) f(j) \quad i, j \geq 1 \quad (2)$$

where $f(k)$ denotes the number of distinct ways of forming a k -mers from k non-distinguishable units with distinguishable functional groups and $1/\lambda$ ($\lambda > 0$) represents the fragmentation strength which, as in [25], can be written as $1/\lambda = \exp(E/KT)$, where E is the Gibbs free energy of a single chemical bond, T is the absolute temperature and K is Boltzmann's constant. Note that we use $1/\lambda$ here instead of λ used by Van Dongen and Ernst in [25] since it is used to call $\lambda > \lambda_c$ (λ_c is a critical value for the occurrence of gelation) as the supercritical stage. Formula (2) and its meaning can also be found in [25]. The choice of $Q_{\underline{n} \underline{n}'}$ reflects the fact that in the homogeneous system (ignoring diffusion effects), reaction occurs with a probability proportional to the number of reactants and inversely proportional to the volume; here the density is taken to be equal to one, so that the volume coincides with the total number of units N .

When $F_{ij} = 0$, the above irreversible random polymerization model was proposed first by Marcus [18] and studied in detail by Lushnikov [17] and Buffet and Pulé [4, 5], which is the stochastic counterpart of Smoluchlovski's coagulation equations, namely the Marcus–Lushnikov coagulation model or process. For readers who are interested in the mathematical aspects of the model, we recommend the survey paper of Aldous [1]. As clusters are growing in size, break-up processes become more important, and the irreversible coagulation reaction should be replaced by coagulation–fragmentation reaction. Van Dongen and Ernst [24, 25] and Spouge [22] were the first ones to extend Smoluchlovski's coagulation equations by including the fragmentation reaction. Motivated by the work done by Van Dongen and Ernst [24, 25], Buffet and Pulé [4, 5] and Pittel *et al* [20, 21], we considered in [13] a reversible random polymerization process (reversible Marcus–Lushnikov process) which is a reversible Markov process with the transition rates $Q_{\underline{n} \underline{n}'}$ with $F_{ij} > 0$ satisfying the detailed balance condition (2). In recent years various aspects of Smoluchlovski's equations and its stochastic counterpart containing the combined effects of coagulation and fragmentation have been extensively studied by many authors (see [2, 6–9, 11–13, 15, 16]).

Although there are many studies devoted to the deterministic and stochastic models based on the coagulation–fragmentation reaction of polymerization, the asymptotic probability distributions of the size of the largest length of the models have received minimal attention. The research work with respect to this problem can be found in the papers by Pittel *et al* [20, 21].

The objective of this paper is to further study the thermodynamic limit distribution of the size of polymerization of our model depending on the fragmentation strength. To explain the meaning of the limit distribution we present some notation and results in the following. From lemmas 1 and 2 of [13] it follows that the process $M_N(t)$ has a unique stationary distribution

$$P_N(\underline{n}) = \frac{1}{\pi_N} \prod_{k=1}^N \left[\left(\frac{N}{\lambda} \right) f(k) \right]^{n_k} / n_k! \quad \underline{n} \in \Omega_N \quad (3)$$

where

$$\pi_N = \pi_N \left(\frac{N}{\lambda} \right) = \sum_{\underline{n} \in \Omega_N} \prod_{k \geq 1} \frac{\left[\frac{N}{\lambda} f(k) \right]^{n_k}}{n_k!} \quad (4)$$

is usually called the partition function of the process. The generating function of the partition functions $\{\pi_N(y)\}$ is given by

$$\sum_{N=0}^{\infty} \pi_N(y) x^N = \exp\{yF(x)\} \quad |x| \leq \bar{r} \quad (5)$$

for any fixed y , where $\pi_0(y) = 1$ and

$$F(x) = \sum_{k=1}^{\infty} f(k)x^k$$

has a positive radius, \bar{r} , of convergence. Furthermore, the partition function has an integral formula

$$\pi_N(y) = \frac{1}{2\pi i} \int_{\Gamma} \exp\{yF(x) - N \log x\} x^{-1} dx$$

where Γ denotes a contour surrounding the origin $x = 0$. In particular,

$$\pi_N = \pi_N \left(\frac{N}{\lambda} \right) = \frac{1}{2\pi i} \int_{\Gamma} \exp \left\{ \frac{N}{\lambda} F(x) - N \log x \right\} x^{-1} dx.$$

Technically, the present paper studies the asymptotic behaviour of the stationary distribution $P_N(\underline{n})$ in the thermodynamic limit as $N \rightarrow \infty$. We shall prove that there exists a critical value λ_c of the fragmentation strength λ such that the size of the largest length of polymers is of order $\log N$, N^a ($\frac{1}{2} < a < 1$), or N , respectively, depending upon whether λ is below λ_c , nearly (or equal to) λ_c , or above λ_c as $N \rightarrow \infty$, where $a = 1/(\beta - 1)$ and β ($2 < \beta < 3$) is an exponent of $f(k)$ in (6).

In section 2 we present the main results on the limit distribution of the k th largest length of polymers in the subcritical, near-critical and supercritical stages and explain the relation of our results to the works done by Pittel *et al* in [20, 21]. The proofs of two theorems are given in section 3. An application of the theorems is shown in section 4.

It should be noted that though the process considered in the paper is different from that studied by Pittel *et al* in [20, 21] and some more general results are obtained in our model, Pittel, Woyczynski and Mann's work provides us with some good ideas and techniques.

2. The main theorems

By theorem 2 of [12] we know that if the positive number $f(k)$ in (2) is of the form $f(k) = (1 + o(1))c\bar{r}^{-k}k^{-\beta}$, where \bar{r} is the positive radius of convergence, c and β are two positive constants, then a necessary and sufficient condition for the occurrence of a gelation in the process is that the number β satisfies $2 < \beta < 3$. So, we shall assume in the following that the positive number $f(k)$ in (2) satisfies

$$f(k) = (1 + o(1))c\bar{r}^{-k}k^{-\beta} \quad (6)$$

where $2 < \beta < 3$. The number β is usually called an exponent of $f(k)$.

It follows from (6) that

$$F'(\bar{r}) = \lim_{x \rightarrow \bar{r}-0} F'(x) < +\infty \quad F''(\bar{r}) = \lim_{x \rightarrow \bar{r}-0} F''(x) = +\infty.$$

According to theorem 2 of [12] a critical value of the fragmentation strength λ for the occurrence of a gelation can be determined as follows:

$$\lambda_c = \bar{r}F'(\bar{r}).$$

It can be checked that the number $f(k)$ for many models, such as RA_a ($a \geq 3$), RA_∞ , A_aRB_b ($\min(a, b) \geq 2$), A_aRB_∞ ($a \geq 2$), etc, satisfies (6), where a and b are two natural numbers (see [13, 26]).

In order to compare the different results corresponding to different domains of λ , let us first recall the main subcritical result, then mention the near-critical and supercritical results.

Let $L_N^{(k)}$ denote the size of k th largest length of polymers and assume that the positive numbers $f(k)$ satisfy (2) and (6) in the following three theorems. We recall first the subcritical result which has been proved in [13]:

Theorem 1. *If $\lambda < \lambda_c$, then the size of the largest length of polymers in \underline{n} (keeping in mind that $\underline{n} = \underline{n}(N)$) is asymptotically (in probability) a logarithmic function of N , i.e.,*

$$L_N^{(1)} = R^{-1}[\log N - \beta \log \log N + O_p(1)]$$

where $R = \log(\bar{r}/r)$, r is the positive root of $\lambda = xF'(x)$, $x \in (0, \bar{r})$, i.e., $\lambda = rF'(r)$, and $O_p(1)$ denotes random variables bounded in probability.

Now we mention the main results of this paper.

Theorem 2. *Suppose that $\lambda = \lambda_c$ or more general, $\lambda = \lambda_N$ satisfies*

$$\lambda_c/\lambda_N = 1 - a(d_\lambda)^{\frac{1}{\beta-1}} N^{-\frac{(\beta-2)}{(\beta-1)}}$$

where $a \in (-\infty, +\infty)$ is fixed and

$$d_\lambda = \frac{c}{\lambda(\beta-1)(\beta-2)(3-\beta)}.$$

Then, for every $x > 0$ and $k \geq 1$,

$$\lim_{N \rightarrow \infty} P(L_N^{(k)} \leq x(d_\lambda N)^{\frac{1}{\beta-1}}) = e^{-I(x)} \sum_{0 \leq j \leq k-1} [I(x)]^j / j!$$

where

$$I(x) = \frac{(\beta-1)(\beta-2)(3-\beta)}{p(a)} \int_x^{+\infty} y^{-\beta} p(a-y) dy$$

and $p(a) = p(a; \beta-1, \delta)$ is the density of a $(\beta-1)$ -stable distribution.

Theorem 3. *Suppose that $\lambda > \lambda_c$. Then*

(i) *The distribution of $L_N^{(1)}$ satisfies a local limit-type relation*

$$P(L_N^{(1)} = j) = (1 + o(1))p(x_j)\Delta x_j$$

where

$$x_j := [N(1 - \lambda_c/\lambda) - j](d_\lambda N)^{-\frac{1}{\beta-1}} \quad \Delta x_j := B(d_\lambda N)^{-\frac{1}{\beta-1}}$$

and

$$B = (3 - \beta)\Gamma(3 - \beta)$$

and $\Gamma(\cdot)$ is the Gamma function. That is, in the distribution,

$$[N(1 - \lambda_c/\lambda) - L_N^{(1)}](d_\lambda N)^{-\frac{1}{\beta-1}} \Longrightarrow X \quad (N \rightarrow +\infty)$$

where the random variable X has the characteristic function

$$E(\exp(itX)) = \exp\{-|t|^{\beta-1} e^{-i\frac{(\beta-3)\pi}{2} \text{sign}(t)}\}.$$

(ii) For every fixed $x > 0$ and $k \geq 2$,

$$\lim_{N \rightarrow \infty} P \left(L_N^{(k)} \leq x N^{\frac{1}{\beta-1}} \right) = e^{-J(x)} \sum_{0 \leq j \leq k-2} [J(x)]^j / j!$$

$$\text{where } J(x) = \frac{c}{\lambda(\beta-1)} x^{-(\beta-1)}.$$

Remark 1. By theorems 1–3 we see that the size of the largest length of polymers is of order $\log N - \beta \log \log N$, $N^{\frac{1}{\beta-1}}$, or N , respectively, depending upon whether λ is below λ_c , nearly (or equal to) λ_c , or above λ_c as $N \rightarrow \infty$.

Note that there is no other restrictive condition on the coagulation rate constants R_{ij} and fragmentation rate constants F_{ij} in the three theorems except for the detailed balance condition. In order to show the relation between the exponent β and R_{ij} and F_{ij} , a form of R_{ij} and F_{ij} will be given in the following.

To model surface interactions, the coagulation and fragmentation coefficients can be taken as

$$R_{ij} = i^\sigma j^\sigma \quad (7)$$

and

$$\sum_{i+j=k} F_{ij} = \frac{2}{\lambda} (k-1)^\sigma \quad (8)$$

where $\sigma \geq 0$. Note that the number $k^\sigma - 1$ proposed by van Dongen and Ernst in [25] is replaced by $(k-1)^\sigma$ here. It follows from proposition 1 in [12] that a necessary and sufficient condition for the occurrence of a gelation is

$$\frac{1}{2} < \sigma < \frac{3}{2} \quad (9)$$

and

$$\sum_{k=1}^{\infty} k^{1+\sigma} f(k) \bar{r}^k = \infty. \quad (10)$$

Moreover, the positive number $f(k)$ has the following form

$$f(k) = (1 + o(1)) c \bar{r}^{-k} k^{-(3/2+\sigma)}$$

if (7), (8), (9) and (10) are given. In this case, $2 < \beta = (3/2 + \sigma) < 3$ and $1/(\beta - 1) = 2/(2\sigma + 1)$. Thus, we have the following corollary.

Corollary 1. Let R_{ij} , F_{ij} and $f(k)$ satisfy (7), (8), (9) and (10) respectively. Then there exists a critical value λ_c such that $\lambda_c = \bar{r} F'(\bar{r}) < \infty$ and the size of the largest length of polymers is of order $\log N - (3/2 + \sigma) \log \log N$, $N^{\frac{2}{2\sigma+1}}$, or N , respectively, depending upon whether λ is below λ_c , nearly (or equal to) λ_c , or above λ_c as $N \rightarrow \infty$.

The model studied by Pittel, Woyczynski and Mann is a Whittle-type random graph process of polymerization [20, 21] which is rather different from the reversible Marcus–Lushnikov process. They have proved in [21] that the size of the largest component is of the order $\log N - \frac{5}{2} \log \log N$, $N^{2/3}$ and N , respectively, in the subcritical, near-critical and supercritical stages. This result is just the special case of corollary 1 for $\sigma = 1$, i.e. $R_{ij} = ij$. In fact, the gelation in the case $R_{ij} = ij$ is known to be equivalent to the emergence of a giant component in the random graph theory, a result which was initiated by Erdős and Rényi [10] and extensively studied by Bollobás [3], Pittel [19] and Janson *et al* [14]. Thus, it is the reversible Marcus–Lushnikov process including the surface interactions ($R_{ij} = i^\sigma j^\sigma$, $1/2 < \sigma < 3/2$) of the polymers that we can obtain more general results.

3. Proofs of the theorems

The key to proving theorems 2 and 3 is to obtain the asymptotic formulae of the partition function π_N in (4).

Let $\lambda_N > 0$, $F_{ij}^{(N)} > 0$ depend on N such that

$$\lambda_N F_{ij}^{(N)} f(i+j) = R_{ij} f(i) f(j) \quad 2 \leq i+j \leq N. \quad (11)$$

If $\lambda_N = \lambda$ is independent of N , then $F_{ij}^{(N)}$ is also, and (11) becomes (2).

Lemma 1. *Let N_k be the random number of k -mers and $E(\cdot)$ denote the expectation corresponding to the stationary distribution $P_N(\cdot)$ in (3). Then*

(i) *For every fixed $k > 0$ and $\lambda = \lambda_c$ or more general, $\lambda = \lambda_N$ satisfying*

$$\lambda_c / \lambda_N = 1 - a_N (d_\lambda)^{\frac{1}{\beta-1}} N^{-\frac{(\beta-2)}{(\beta-1)}}$$

we have

$$E(N_k) = (1 + o(1)) c \lambda^{-1} k^{-\beta} N p(a_N - k (d_\lambda N)^{-\frac{1}{\beta-1}}) / p(a_N)$$

where a_N is uniformly bounded to N .

(ii) *For $k < l_N$ and $\lambda > \lambda_c$,*

$$E(N_k) = (1 + o(1)) c \lambda^{-1} b^\beta [k(b - k/N)]^{-\beta} N.$$

(iii) *For $l_N < k < L_N$ and $\lambda > \lambda_c$,*

$$E(N_k) = (1 + o(1)) B b^\beta (k/N)^{-\beta} (d_\lambda N)^{-\frac{1}{\beta-1}} p(x_{N,k}).$$

(iv) *For $k > L_N$ and $\lambda > \lambda_c$,*

$$E(N_k) = (1 + o(1)) c \lambda^{-1} b^\beta [k(k/N - b)]^{-\beta} N$$

where $x_{N,k} = (b - k/N)(d_\lambda)^{-\frac{1}{\beta-1}} N^{\frac{(\beta-2)}{(\beta-1)}}$, $b = 1 - \frac{\lambda_c}{\lambda}$, $l_N = bN - \omega(N)N^{1/(\beta-1)}$, $L_N = bN + \omega(N)N^{1/(\beta-1)}$ and $\omega(N) \rightarrow +\infty$ however slowly.

Proof of lemma 1. It follows from (3), (4) and (6) that

$$\begin{aligned} E(N_k) &= \sum_{\underline{n} \in \Omega_N} n_k P_N(\underline{n}) \\ &= \frac{Nf(k)}{\lambda \pi_N} \sum_{\underline{n} \in \Omega_N} \frac{[\frac{N}{\lambda} f(k)]^{n_k-1}}{(n_k-1)!} \prod_{j \neq k} \frac{[\frac{N}{\lambda} f(j)]^{n_j}}{n_j!} \\ &= \frac{Nf(k)}{\lambda \pi_N} \sum_{\underline{n} \in \Omega_{N-k}} \prod_{j=1}^{N-k} \frac{[\frac{N}{\lambda} f(j)]^{n_j}}{n_j!} \\ &= \frac{Nf(k)}{\lambda} \frac{\pi_{N-k}(\frac{N}{\lambda})}{\pi_N(\frac{N}{\lambda})} = (1 + o(1)) c \frac{N}{\lambda} \bar{r}^{-k} k^{-\beta} \frac{\pi_{N-k}(\frac{N}{\lambda})}{\pi_N(\frac{N}{\lambda})}. \end{aligned} \quad (12)$$

Note that the numbers, $\pi_N(\frac{N}{\lambda})$ and $\pi_{N-k}(\frac{N}{\lambda})$, have been estimated in the proof of theorem 2 of [12]. That is

(I) If $\lambda = \lambda_c$, or more generally, $\lambda = \lambda_N$ satisfies $\lambda_c/\lambda_N = 1 - a_N(d_\lambda)^{\frac{1}{\beta-1}} N^{-\frac{(\beta-2)}{(\beta-1)}}$, where a_N is uniformly bounded to N , then

$$\pi_N \left(\frac{N}{\lambda} \right) = (1 + o(1)) D_N(\bar{r}) p(a_N) (d_\lambda N)^{-\frac{1}{\beta-1}}$$

and, for every $k > 0$,

$$\pi_{N-k} \left(\frac{N}{\lambda} \right) = (1 + o(1)) (\bar{r})^k D_N(\bar{r}) p(a_N - k(d_\lambda N)^{-\frac{1}{\beta-1}}) (d_\lambda N)^{-\frac{1}{\beta-1}}$$

where $D_N(\bar{r}) = \exp \left\{ \frac{N}{\lambda} F(\bar{r}) - N \log \bar{r} \right\}$.

(II) If $\lambda > \lambda_c$, then

$$\pi_N \left(\frac{N}{\lambda} \right) = (1 + o(1)) d_\lambda \left(1 - \frac{\lambda_c}{\lambda} \right)^{-\beta} D_N(\bar{r}) p \left(0; \frac{1}{\beta-1}, 2 - \frac{3}{\beta-1} \right) N^{-(\beta-1)} \quad (13)$$

and

$$\pi_{N-k} \left(\frac{N}{\lambda} \right) = (1 + o(1)) d_\lambda (\bar{r})^k (b - k/N)^{-\beta} D_N(\bar{r}) p \left(0; \frac{1}{\beta-1}, 2 - \frac{3}{\beta-1} \right) N^{-(\beta-1)}$$

for $k < l_N$,

$$\pi_{N-k} \left(\frac{N}{\lambda} \right) = (1 + o(1)) (\bar{r})^k D_N(\bar{r}) p(x_{N,k}) (d_\lambda N)^{-\frac{1}{\beta-1}}$$

for $l_N < k < L_N$, and

$$\begin{aligned} \pi_{N-k} \left(\frac{N}{\lambda} \right) &= (1 + o(1)) d_\lambda (\bar{r})^k (k/N - b)^{-\beta} D_N(\bar{r}) \\ &\quad \times p \left(0; \frac{1}{\beta-1}, - \left(2 - \frac{3}{\beta-1} \right) \right) N^{-(\beta-1)} \end{aligned}$$

for $k > L_N$. Note that $\Gamma(s+1) = s\Gamma(s)$, $\Gamma(s)\Gamma(1-s) = \pi/\sin(s\pi)$ for $0 < s < 1$, $p(x; \beta-1, \delta) = p(-x; \beta-1, -\delta)$ and

$$p \left(0; \frac{1}{\beta-1}, 2 - \frac{3}{\beta-1} \right) = \pi^{-1} \Gamma(\beta) \sin[(\beta-2)\pi]$$

(see [23]). Thus, by (6), (12), (I) and (II), we can immediately obtain (i), (ii), (iii) and (iv) of lemma 1. \square

Proof of theorem 2. We first prove that $L_N^{(1)}/N^{\frac{1}{\beta-1}}$ is bounded in probability. That is we will prove that there is no polymer of size $k \geq k_N = \omega(N)(d_\lambda N)^{\frac{1}{\beta-1}}$, where $\omega(N) \rightarrow +\infty$ however slowly. By (i) of lemma 1 we have

$$\begin{aligned} E \left(\sum_{k \geq k_N} N_k \right) &= \sum_{k \geq k_N} E(N_k) \\ &= O \left(\sum_{k \geq k_N} \left(\frac{k}{N^{\frac{1}{\beta-1}}} \right)^{-\beta} p(a_N - k(d_\lambda N)^{-\frac{1}{\beta-1}}) \frac{1}{N^{\frac{1}{\beta-1}}} \right) \\ &= O \left(\int_{\omega(N)}^{(d_\lambda)^{-\frac{1}{\beta-1}} N^{\frac{\beta-2}{\beta-1}}} x^{-\beta} p(a_N - x) dx \right) \\ &= O \left((\omega(N))^{-(\beta-1)} \right) = o(1). \end{aligned}$$

For $y > x$ we set $k_1 = x(d_\lambda N)^{\frac{1}{\beta-1}}$, $k_2 = y(d_\lambda N)^{\frac{1}{\beta-1}}$, $C_N(x, y) = \sum_{k=k_1}^{k_2} N_k$ and

$$I(x, y) = \frac{(\beta-1)(\beta-2)(3-\beta)}{p(a)} \int_x^y u^{-\beta} p(a-u) du.$$

Let $a_N \rightarrow a$ as $N \rightarrow \infty$. It follows from (i) of lemma 1 that

$$\begin{aligned} E(C_N(x, y)) &= (1+o(1)) \frac{(\beta-1)(\beta-2)(3-\beta)}{p(a_N)} \sum_{k_1 \leq k \leq k_2} \left(\frac{k}{(d_\lambda N)^{\frac{1}{\beta-1}}} \right)^{-\beta} \\ &\quad \times p \left(a_N - k(d_\lambda N)^{-\frac{1}{\beta-1}} \right) \frac{1}{(d_\lambda N)^{\frac{1}{\beta-1}}} \\ &= (1+o(1)) \frac{(\beta-1)(\beta-2)(3-\beta)}{p(a_N)} \int_x^y u^{-\beta} p(a_N - u) du \\ &\rightarrow I(x, y) \end{aligned}$$

as $N \rightarrow \infty$. Furthermore,

$$\begin{aligned} E(C_N(x, y)(C_N(x, y) - 1)) &= \sum_{k_1 \leq k \leq k_2} \sum_{k_1 \leq l \leq k_2} \frac{N^2 f(k)f(l)}{\lambda^2} \frac{\pi_{N-k-l} \left(\frac{N}{\lambda} \right)}{\pi_N \left(\frac{N}{\lambda} \right)} \\ &= (1+o(1)) \sum_{k_1 \leq k \leq k_2} \sum_{k_1 \leq l \leq k_2} \left(\frac{cN}{\lambda} \right)^2 (kl)^{-\beta} p(a_N - (k+l)(d_\lambda N)^{-\frac{1}{\beta-1}}) / p(a_N) \\ &= (1+o(1)) \frac{[(\beta-1)(\beta-2)(3-\beta)]^2}{p(a_N)^2} \\ &\quad \times \int_x^y \int_x^y u^{-\beta} v^{-\beta} p(a_N) p(a_N - u - v) du dv \\ &\rightarrow [I(x, y)]^2 \end{aligned}$$

as $N \rightarrow \infty$. By the same method we have, for every $k \geq 1$,

$$E([C_N(x, y)]_k) \rightarrow [I(x, y)]^k$$

as $N \rightarrow \infty$, where $[C_N(x, y)]_k = C_N(x, y)(C_N(x, y) - 1)(C_N(x, y) - 2) \dots (C_N(x, y) - k + 1)$ is the total number of ordered k -tuples of different polymers of size $j \in [k_1, k_2]$. Hence, $C_N(x, y)$ is, in the limit, a Poisson distribution with parameter $I(x, y)$. Since $L_N^{(1)} / N^{\frac{1}{\beta-1}} = O_p(1)$, it follows that, for every $x > 0$ $k \geq 1$ and large x_1 ,

$$\begin{aligned} P(L_N^{(k)} \geq x(d_\lambda N)^{\frac{1}{\beta-1}}) &= P(x(d_\lambda N)^{\frac{1}{\beta-1}} \leq L_N^{(k)} \leq L_N^{(1)} \leq x_1(d_\lambda N)^{\frac{1}{\beta-1}}) + O(P(L_N^{(1)} > x_1(d_\lambda N)^{\frac{1}{\beta-1}})) \\ &= P(C_N(x, y) \geq k) + O(P(L_N^{(1)} > x_1(d_\lambda N)^{\frac{1}{\beta-1}})) \rightarrow e^{-I(x)} \sum_{j \geq k} [I(x)]^j / j! \end{aligned}$$

as $x_1 \rightarrow \infty$ and $N \rightarrow \infty$, where $I(x) = \lim_{y \rightarrow \infty} I(x, y)$. This completes the proof of theorem 2. \square

Proof of theorem 3. (i) By (iii) of lemma 1 we have

$$\begin{aligned} E(N_j) &= (1+o(1)) B b^\beta (j/N)^{-\beta} (d_\lambda N)^{-\frac{1}{\beta-1}} p(x_{N,j}) \\ &= (1+o(1)) p(x_j) \Delta x_j \end{aligned} \tag{14}$$

where $j = N(1 - \lambda_c/\lambda) - x_j(d_\lambda N)^{\frac{1}{\beta-1}}$ and $\Delta x_j = B(d_\lambda N)^{-\frac{1}{\beta-1}}$. Set $C_N(j) = \sum_{i \geq j} N_i$. Then

$$E(N_j) = P(N_j = 1) + \sum_{k \geq 2} k P(N_j = k) = P(N_j = 1) + O[NP(C_N(j) \geq 2)]$$

and

$$\begin{aligned} P(N_j = 1) &= P(N_j = 1, L_N^{(1)} = j) + P(N_j = 1, L_N^{(1)} > j) \\ &= P(L_N^{(1)} = j) - P(N_j \geq 2, L_N^{(1)} = j) + P(N_j = 1, L_N^{(1)} > j) \\ &= P(L_N^{(1)} = j) + O[P(C_N(j) \geq 2)]. \end{aligned}$$

Therefore

$$P(L_N^{(1)} = j) = E(N_j) + O[NP(C_N(j) \geq 2)]. \quad (15)$$

Furthermore, for $j_1, j_2 \geq j$,

$$N' = N - (j_1 + j_2) = N(2\lambda_c/\lambda - 1) + O(N^{\frac{1}{\beta-1}}).$$

If $\lambda_c/\lambda < 1/2$, then $N' < 0$. This means that $P(C_N(j) \geq 2) = 0$. Let $\lambda_c/\lambda \geq 1/2$ and $\lambda_{N'} = \lambda N'/N$, we have

$$\lambda_{N'} \leq 2\lambda_c - \lambda + o(1) \leq \lambda_c - (\lambda - \lambda_c)/2 < \lambda_c.$$

Let $r' \in (0, \bar{r})$ satisfying $\lambda_{N'} = r' F'(r')$. It follows from (5) that

$$\pi_{N'} \left(\frac{N}{\lambda} \right) \leq \exp \left\{ \frac{N}{\lambda} F(x) - N' \log x \right\}$$

for any $x > 0$. So,

$$\pi_{N'} \left(\frac{N}{\lambda} \right) \leq \exp \left\{ \frac{N}{\lambda} F(r') - N' \log r' \right\}. \quad (16)$$

Thus, by (6), (13) and (16), we have

$$\begin{aligned} P(C_N(j) \geq 2) &\leq E([C_N(j)]_2) \\ &= \sum_{j_1, j_2 \geq j} \left(\frac{N}{\lambda} \right)^2 f(j_1) f(j_2) \frac{\pi_{N-(j_1+j_2)} \left(\frac{N}{\lambda} \right)}{\pi_N \left(\frac{N}{\lambda} \right)} \\ &\leq O \left(\sum_{j_1, j_2 \geq j} (j_1 j_2)^{-\beta} N^{\beta+1} \exp \left\{ \frac{N}{\lambda} (F(r') - F(\bar{r})) - N' (\log r' - \log \bar{r}) \right\} \right) \\ &= O \left(N^{3-\beta} \exp \left\{ \frac{N}{\lambda} (F(r') - F(\bar{r})) - N' (\log r' - \log \bar{r}) \right\} \right). \end{aligned}$$

Since

$$\begin{aligned} &\frac{N}{\lambda} (F(r') - F(\bar{r})) - N' (\log r' - \log \bar{r}) \\ &= \frac{N}{\lambda} \left[- \int_{r'}^{\bar{r}} x F'(x) x^{-1} dx + r' F'(r') \int_{r'}^{\bar{r}} x^{-1} dx \right] \\ &= - \frac{N}{\lambda} \int_{r'}^{\bar{r}} [x F'(x) - r' F'(r')] x^{-1} dx \\ &= - \frac{N}{\lambda} c' \end{aligned}$$

where $c' > 0$, it follows that

$$NP(C_N(j) \geq 2) \leq O\left(N^{4-\beta} \exp\left\{-\frac{N}{\lambda}c'\right\}\right) = o(N^{-\frac{1}{\beta-1}}).$$

Hence, by (14) and (15),

$$\begin{aligned} P(L_N^{(1)} = j) &= E(N_j) + o\left(N^{-\frac{1}{\beta-1}}\right) \\ &= (1 + o(1))p(x_j)\Delta x_j. \end{aligned}$$

This completes the proof of (i).

(ii) We first prove that there is at most one polymer of size $j \geq j_N = \omega(N)N^{\frac{1}{\beta-1}}$, where $\omega(N) \rightarrow +\infty$ however slowly. For any $0 < \epsilon < b = 1 - \frac{\lambda_c}{\lambda}$ and $\epsilon + b \leq 1$, let $l_N(\epsilon) = bN - \epsilon N$. It follows from (ii) of lemma 1 that

$$\begin{aligned} E\left(\sum_{j_N \leq j \leq l_N(\epsilon)} N_j\right) &= (1 + o(1)) \sum_{j_N \leq j \leq l_N(\epsilon)} c\lambda^{-1}b^\beta [j(b - j/N)]^{-\beta} N \\ &= O\left(N^{-(\beta-2)} \int_{j_N/N}^{b-\epsilon} x^{-\beta}(b-x)^{-\beta} dx\right) \\ &= O((\omega(N))^{-(\beta-1)}) \rightarrow 0 \end{aligned}$$

as $N \rightarrow \infty$. Therefore, there is no polymer in the ranges from j_N to $l_N(\epsilon)$. We now consider the range from $l_N(\epsilon)$ to N . Suppose that there are two polymers j_1 and j_2 of size between $l_N(\epsilon)$ and N . If $\frac{\lambda_c}{\lambda} < 1/2$, then

$$N - (j_1 + j_2) \leq N - 2l_N(\epsilon) = N \left(2\frac{\lambda_c}{\lambda} - 1 + 2\epsilon\right) < 0$$

for all small enough $\epsilon > 0$. This means that there is at most one polymer of size between $l_N(\epsilon)$ and N for $\frac{\lambda_c}{\lambda} < 1/2$. Let $\frac{\lambda_c}{\lambda} \geq 1/2$. Then, taking $\epsilon < b/2$ we have

$$0 \leq N' = N - (j_1 + j_2) \leq N - 2l_N(\epsilon) = N \left(2\frac{\lambda_c}{\lambda} - 1 + 2\epsilon\right)$$

and

$$\lambda_{N'} \leq \lambda \left(2\frac{\lambda_c}{\lambda} - 1 + 2\epsilon\right) < \lambda_c$$

where $\lambda_{N'} = \lambda N'/N$. So, as in the proof of (i),

$$E\left(\sum_{l_N(\epsilon) \leq j_1, j_2 \leq N} N_j\right) \leq E([C_N(l_N(\epsilon))]_2) \leq O\left(N^{3-\beta} \exp\left\{-\frac{N}{\lambda}c''\right\}\right) \rightarrow 0$$

as $N \rightarrow \infty$, where $c'' > 0$. Thus, we have proved that there is at most one polymer of size $j \geq \omega(N)N^{\frac{1}{\beta-1}}$ as $N \rightarrow \infty$. Therefore, for each $\epsilon > 0$, there exists $x_1 = x_1(\epsilon) > x$ so large that

$$\limsup P(L_N^{(2)} > x_1 N^{\frac{1}{\beta-1}}) < \epsilon.$$

Note that, for every $k \geq 2$,

$$\begin{aligned} P(L_N^{(k)} \geq xN^{\frac{1}{\beta-1}}) &= P(xN^{\frac{1}{\beta-1}} \leq L_N^{(k)} \leq L_N^{(2)} \leq x_1 N^{\frac{1}{\beta-1}}) + O(P(L_N^{(2)} > x_1 N^{\frac{1}{\beta-1}})) \\ &= P(C_N(x, x_1) \geq k - 1) + O(P(L_N^{(2)} > x_1 N^{\frac{1}{\beta-1}})) \end{aligned}$$

where $C_N(x, x_1) = \sum_{j_1 \leq j \leq j_2} N_j$, $j_1 = xN^{\frac{1}{\beta-1}}$ and $j_2 = x_1N^{\frac{1}{\beta-1}}$. So, as in the proof of theorem 2, $C_N(x, x_1)$ is, in the limit, a Poisson distribution with parameter

$$J(x, x_1) = \frac{c}{\lambda(\beta-1)} [x^{-(\beta-1)} - x_1^{-(\beta-1)}].$$

Thus

$$\lim_{N \rightarrow \infty} \sup \left| P(L_N^{(k)} > xN^{\frac{1}{\beta-1}}) - e^{-J(x, x_1)} \sum_{j \geq k-1} [J(x, x_1)]^j / j! \right| \leq \epsilon$$

and, letting $\epsilon \rightarrow 0$, i.e. $x_1(\epsilon) \rightarrow +\infty$, we have

$$\lim_{N \rightarrow \infty} P(L_N^{(k)} \leq xN^{\frac{1}{\beta-1}}) = e^{-J(x)} \sum_{0 \leq j \leq k-2} [J(x)]^j / j!$$

for every $x > 0$ and $k \geq 2$, where $\lim_{x_1 \rightarrow +\infty} J(x, x_1) = J(x) = \frac{c}{\lambda(\beta-1)} x^{-(\beta-1)}$. This completes the proof. \square

4. Applications

As an application of theorems 1–3, we show two examples.

Example 1. RA_a model ($a \geq 3$).

It is known that the numbers $f(k)$ for the RA_a model are as follows:

$$f(k) = \frac{a^k [(a-1)k!]}{k! [(a-2)k+2]}.$$

By using Stirling's formula we have (see [13])

$$f(k) = (1 + o(1)) c \bar{r}^{-k} k^{-5/2}$$

where $c = \sqrt{(a-1)/[2\pi(a-2)^5]}$, $\beta = 5/2$

$$\bar{r} = \lim_{k \rightarrow \infty} \frac{f(k)}{f(k+1)} = \frac{(a-2)^{(a-2)}}{a(a-1)^{(a-1)}}$$

and $\lambda_c = \bar{r}F'(\bar{r}) = (a-1)/[a(a-2)^2]$. Taking the fragmentation coefficients $F(i, j)$ such that

$$F_{ij} = \frac{1}{\lambda} \frac{R_{ij} f(i) f(j)}{f(i+j)} \quad i, j \geq 1$$

it follows from theorems 1–3 that the size of the largest length of polymers is of order $\log N - \frac{5}{2} \log \log N$, $N^{2/3}$ or N , respectively, depending upon whether λ is below λ_c , nearly (or equal to) λ_c , or above λ_c as $N \rightarrow \infty$.

Example 2. For the model RA_∞ we have

$$f(k) = \frac{k^{k-2}}{k!}.$$

It can be calculated that

$$f(k) = (1 + o(1)) c \bar{r}^{-k} k^{-5/2}$$

where $c = (2\pi)^{-1/2}$ and $\bar{r} = e^{-1}$. Therefore, the size of the largest length of polymers is of order $\log N - \frac{5}{2} \log \log N$, $N^{2/3}$ and N , respectively, in the subcritical, near-critical and the supercritical cases.

Acknowledgment

I want to thank the referees for their valuable suggestions. This work is partly supported by National Science Foundation of China.

References

- [1] Aldous D J 1999 Deterministic and stochastic models for coalescence (aggregation and coagulation): a review of the mean-field theory for probabilists *Bernoulli* **5** 3–48
- [2] Ball J M and Carr J 1990 The discrete coagulation–fragmentation equation: existence, uniqueness, and density conservation *J. Stat. Phys.* **61** 203–34
- [3] Bollobás B 1985 *Random Graphs* (London: Academic)
- [4] Buffet E and Pulé J V 1990 On the Lushnikov’s model of gelation *J. Stat. Phys.* **58** 1041–58
- [5] Buffet E and Pulé J V 1991 Polymers and random graphs *J. Stat. Phys.* **64** 87–110
- [6] Dubovskii P B and Stewart I W 1996 Existence, uniqueness and mass conservation for the coagulation–fragmentation equation *Math. Methods Appl. Sci.* **19** 571–91
- [7] Durrett R, Granovsky B and Gueron S 1999 The equilibrium behaviour of reversible coagulation–fragmentation processes *J. Theor. Probab.* **12** 447–74
- [8] Eibeck A and Wagner W 2000 Approximative solution of the coagulation–fragmentation equation by stochastic particle systems *Stoch. Anal. Appl.* **18** 921–48
- [9] Freiman G and Granovsky B 2002 Asymptotic formula for a partition function of reversible coagulation–fragmentation processes *J. Isr. Math.* **130** 259–79
- [10] Erdős P and Rényi A 1960 On the evolution of random graphs *Magy. Tud. Akad. Mat. Kut. Intéz. Közl.* **5** 17–61
- [11] Guias F 2001 Convergence properties of a stochastic model for coagulation–fragmentation processes with diffusion *Stoch. Anal. Appl.* **19** 245–78
- [12] Han D 2003 A necessary and sufficient condition for gelation of a reversible Markov process of polymerization *J. Phys. A: Math. Gen.* **36** 893–909
- [13] Han D 1995 Subcritical asymptotic behaviour in the thermodynamic limit of reversible random polymerization process *J. Stat. Phys.* **80** 389–404
- [14] Janson S, Knuth D E, Luczak T and Pittel B 1993 The birth of giant component *Random Struct. Algorithms* **4** 233–58
- [15] Jeon I 1998 Existence of gelling solutions for coagulation fragmentation equations *Commun. Math. Phys.* **194** 541–67
- [16] Laurenzi I J and Diamond S L 2003 Kinetics of random aggregation–fragmentation processes with multiple components *Phys. Rev. E* **67** 051103(1–15)
- [17] Lushnikov A A 1978 Certain new aspects of the coagulation theory *Izv. Atm. Ok. Fiz.* **14** 738–43
- [18] Marcus A H 1968 Stochastic coalescence *Technometrics* **10** 133–43
- [19] Pittel B 1990 On tree census and the giant component in sparse random graphs *Random Struct. Algorithms* **1** 311–42
- [20] Pittel B, Wocczynski W A and Mann J A 1987 From Gaussian subcritical to Holtsmark (3/2-lévy stable) supercritical asymptotic behaviour in ‘rings forbidden’ Flory–Stockmayer model of polymerization *Graph Theory and Topology in Chemistry* ed R B King and D H Rouvray (Amsterdam: Elsevier) pp 362–70
- [21] Pittel B, Woczynski W A and Mann J A 1990 Random tree-type partitions as a model for acyclic polymerization: Holtsmark ($\frac{3}{2}$ -stable) distribution of the supercritical gel *Ann. Probab.* **18** 319–41
- [22] Spouge J L 1984 An existence theorem for the discrete coagulation–fragmentation equations *Math. Proc. Camb. Phil. Soc.* **96** 351–7
- [23] Uchaikin V V and Zolotarev V M 1999 Chance and stability stable distributions and their applications *Modern Probability and Statistics* (Leiden: VSP Intl Science)
- [24] Van Dongen P G J and Ernst M H 1983 Pre- and postgel size distributions in (ir)reversible polymerization *J. Phys. A: Math. Gen.* **16** L327–32
- [25] Van Dongen P G J and Ernst M H 1984a Kinetics of reversible polymerization *J. Stat. Phys.* **37** 301–24
- [26] Van Dongen P G J and Ernst M H 1984b Size distribution in the polymerization model A_fRB_g *J. Phys. A: Math. Gen.* **17** 2281–97